## The onset of chaos: Full renormalization-group calculation

(Sethna, "Entropy, Order Parameters, and Complexity", ex. 12.XXX) © 2017, James P. Sethna, all rights reserved.

In this exercise, we implement Feigenbaum's numerical scheme for finding high-precision values of the universal constants

 $\alpha = -2.50290787509589282228390287322$ 

 $\delta = 4.66920160910299067185320382158$ ,

that quantify the scaling properties of the period-doubling route to chaos (Fig. 12.17}, Exercise 'Period doubling'). This extends the lowest-order calculation of the companion Exercise 'The onset of chaos: Lowest order renormalization-group for period doubling'}.

αFeigenbaum = -2.502907875095892822283902873218; δFeigenbaum = 4.669201609102990671853203821578;

Our renormalization group operation (Exercises 'Period doubling and the renormalization group' and the companion Exercise) coarse-grains in time taking  $g \rightarrow g \circ g$ , and then rescales distance x by a factor of  $\alpha$ . Centering our functions at x=0, this leads to

 $T[g](x) = \alpha g(g(x/\alpha)).$ 

We shall solve for the properties at the onset of chaos by analyzing our function-space renormalizationgroup by expanding our functions in a power series

 $g(\mathbf{x})\approx 1\,+\, \boldsymbol{\Sigma}^{N}{}_{n=1}\,\boldsymbol{G}_{n}\,\boldsymbol{x}^{2\,n}.$ 

Notice that we only keep even powers of x; the fixed point is known to be symmetric about the maximum, and the unstable mode responsible for the exponent  $\delta$  will also be symmetric.

```
(* N us a reserved variable; use Nn instead *)

g[Nn_{1}][x_{1}] := 1 + Sum[..., \{n, 1, Nn\}]

T[g_{1}][x_{1}] := \alpha ...

(* We'll also want the derivative of g later *)

Dg[Nn_{1}][x_{1}] := Sum[..., \{...\}]

(* Test your functions by plotting

them.G=[-1.5,0,0,...] should give T[g] close to g *)

G0[1] := -3/2

G0[n_{1}] := 0 /; n \neq 1

Plot[\{g[2][x], T[g[2]][x] /. \{\alpha \rightarrow 1/g[2][1]\}\} /. G \rightarrow G0, \{x, 0, 2\}]
```

First, we must approximate the fixed point  $g^*(x)$  and the corresponding value of the universal constant  $\alpha$ . At order N, we must solve for  $\alpha$  and the N polynomial coefficients  $G^*_n$ . We can use the N+1 equations fixing the function at equally spaced points in the positive unit interval:

 $T[g \star](x_m)=g^{\star}(x_m), \qquad x_m=m/N, m=\{0,...,N\}.$ We can use the first of these equations to solve for  $\alpha$ . (a) Show that the equation for m=0 sets  $\alpha = 1/g^{*}(1)$ .

We can use a root-finding routine to solve for  $G^*_n$ .

(b) Implement the other N constraint equations above in a form appropriate for your method of finding roots of nonlinear equations, substituting your value for  $\alpha$  from part (a). Check that your routine at N=1 gives values for  $\alpha \approx -2.5$  and G $*1 \approx -1.5$ . (These should reproduce the values from the companion Exercise part (c).)

```
Nmax = 20;

For [Nn = 1, Nn <= Nmax, Nn = Nn + 1,

xm = Range[...];

vars = Table[{G[n], G0[n]}, {...}];

eqns = Table[\ldots = \ldots /. \alpha \rightarrow 1/g[Nn][1], \{x, xm\}];

GStar[Nn] = FindRoot[\ldots, WorkingPrecision \rightarrow 50];

\alpha[Nn] = \ldots] /. GStar[Nn]]

Table[{Nn, \ldots}, {Nn, 1, Nmax}] // MatrixForm

Table[{Nn, \alphaFeigenbaum - \ldots
```

Now we need to solve for the renormalization group flows T[g], linearized about the fixed point  $g(x)=g^*(x)+\epsilon\psi(x)$ . Feigenbaum tells us that  $T[g^*+\epsilon\psi]=T[g_*]+\epsilon\mathcal{L}[\psi]$ , where  $\mathcal{L}$  is the linear operator taking  $\psi(x)$  to

$$\mathcal{L}[\psi](\mathsf{x}) = \alpha \ \psi(\ \mathsf{g}^*(\mathsf{x}/\alpha)) + \alpha \ \mathsf{g}^{*'}(\mathsf{g}(\mathsf{x}/\alpha)) \ \psi(\mathsf{x}/\alpha).$$

(d) Derive the equation above.

## ANSWER HERE

We want to find eigenfunctions that satisfy  $\mathcal{L}[\psi] = \lambda \psi$ . Again, we can expand  $\psi(x)$  in a polynomial  $\psi(x) = \sum_{n=0}^{N-1} \psi_n x^{2n}$  ( $\psi_0 \equiv 1$ ).

We then approximate the action of  $\mathcal{L}$  on  $\psi$  by its action at N points x<sub>i</sub>, that need not be the same as the N points x<sub>m</sub> we used to find g\*. We shall use x<sub>i</sub> = (i-1)/(N-1), i=1,...,N. (For N=1, we use x<sub>1</sub>=0.) This leads us to a linear system of N equations for the coefficients  $\psi_n$ , using the previous two equations.  $\sum^{N-1}_{n=0} [\alpha g(x_i/\alpha)^{2n} + \alpha g'(g(x_i/\alpha)) (x_i/\alpha)^{2n}] \psi_n = \lambda \sum^{N-1}_{n=0} x^{2n}_i \psi_n$ 

These equations for the coefficients  $\psi_n$  of the eigenfunctions of  $\mathbb{B}$  is in the form of a generalized eigenvalue problem

 $\sum_{n=0}^{N-1} L_{\text{in}} \psi_n = \lambda \sum_{n=0}^{N-1} X_{\text{in}} \psi_n.$ 

The solution to the generalized eigenvalue problem can be found from the eigenvalues of  $X^{-1}L$ , but most eigenvalue routines provide a more efficient and accurate option for directly solving the generalized equation given L and X.

(e) Write a routine that calculates the matrices L and X implicitly defined by the previous two equations. For N=1 you should generate 1×1 matrices. For N=1, what is your prediction for  $\delta$ ? (These should reproduce the values from the companion Exercise part (d).)

```
(* Make sure your matrix hasn't transposed rows (i)
 and columns (n).Each row should give powers of one x_i. *)
(* Avoid 0^0 for N=1 by using 'Evaluate' to set up n=0 column in X *)
xtildes[Nn_] := Range[0, 1, 1/(Nn - 1)]
xtildes[1] := {0.}
X[Nn_] := Table[Evaluate[Table[..., {n, 0, Nn - 1}]], {xtilde, xtildes[Nn]}]
X[1]
X[3]
L[Nn_] :=
 Table[Evaluate[Table[\alpha[Nn] g[...]^(...) + \alpha[Nn] Dg[...] (...)^(...) /. GStar[Nn],
    {n, 0, Nn - 1}]], {xtilde, xtildes[Nn]}]
L[1]
L[3] // MatrixForm
Eigenvalues[{L[3], X[3]}]
Nmax = 20;
For[Nn = 1, Nn <= Nmax, Nn = Nn + 1,</pre>
 eigvals = ...[{L[Nn], X[Nn]}];
 \delta[Nn] = eigvals[[1]]]
Table[...] // ...
Table[...] // ...
```