

(Add) K, ν, σ, μ

P218 F01
Lecture 18

①

Tensors of Elasticity

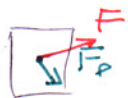
The Stress Tensor


Consider a small planar region $\Delta x \Delta z$ passing through a point \vec{r} . What is the force of the material to the left of \vec{r} on the material to the right?

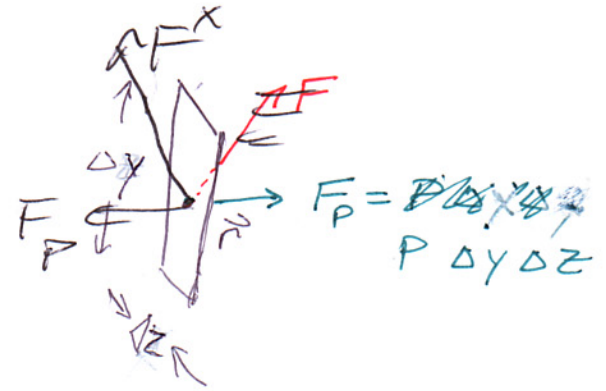
Proportional to $\Delta y \Delta z$:
divide out

Define

$$\sigma_{i1} = F_i / \Delta y \Delta z$$


$$\sigma_{i2} = F_i / \Delta z \Delta x$$


$$\sigma_{i3} = F_i / \Delta x \Delta y$$



Exercise: If the material is a fluid, at pressure P , what is the force?

$$\vec{F} = P \Delta y \Delta z \hat{x}$$

$$\frac{F_1}{\Delta y \Delta z} = \sigma_{11} = P$$

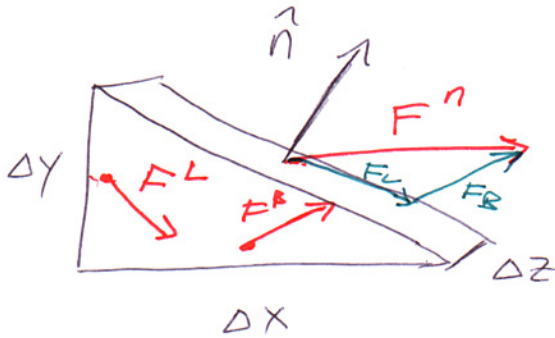
$$F_2 = F_3 = 0$$

σ_{ij} = Stress Tensor

$$\frac{1}{3} \sigma_{ii} = (\sigma_{xx} + \sigma_{yy} + \sigma_{zz}) / 3 = \text{Pressure}$$

Fact $\sigma_{ij} \hat{n}_j = \left(\frac{\text{Force}}{\text{Area}} \right)$ across surface perpendicular to \hat{n}

Proof in Two Dimensions:



Assume tiny triangle of material in equilibrium: force

$$\vec{F}^L + \vec{F}^B = \vec{F}^n$$

(As size $\rightarrow 0$, mass $\rightarrow 0$, acceleration finite \rightarrow net force = 0)

$$\vec{F}_i^n = \sigma_{i1} \Delta y \Delta z + \sigma_{i2} \Delta x \Delta z$$

$$\hat{n} = \frac{1}{\sqrt{\Delta x^2 + \Delta y^2}} \begin{pmatrix} \Delta y \\ \Delta x \end{pmatrix}$$

$$\vec{F}_i^L = \Delta y \Delta z \sigma_{i1}$$

$$\vec{F}_i^B = \Delta x \Delta z \sigma_{i2}$$

$$\text{Area of Face} = \Delta z \sqrt{\Delta x^2 + \Delta y^2}$$

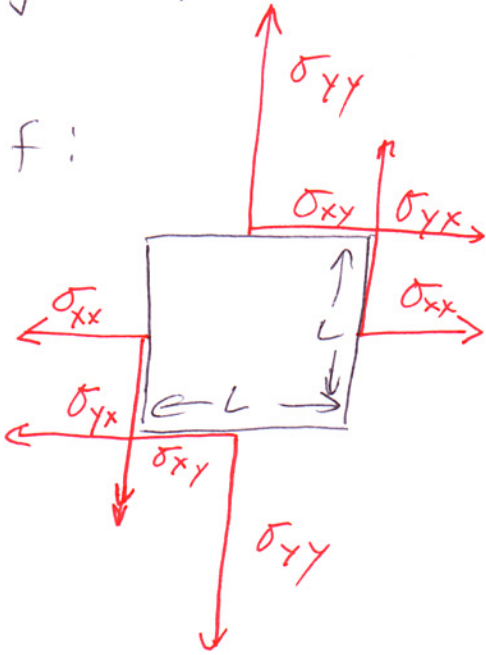
$$\frac{\vec{F}_i^n}{\text{Area}} = \sigma_{i1} \frac{\Delta y}{\sqrt{\Delta x^2 + \Delta y^2}} + \sigma_{i2} \frac{\Delta x}{\sqrt{\Delta x^2 + \Delta y^2}}$$

$$= \sigma_{i1} \hat{n}_1 + \sigma_{i2} \hat{n}_2 = \sigma_{ij} \hat{n}_j \quad \checkmark$$

Fact

$$\sigma_{ij} = \sigma_{ji} \quad (\text{Stress is a symmetric tensor})$$

Proof:



Assume cube is
in equilibrium

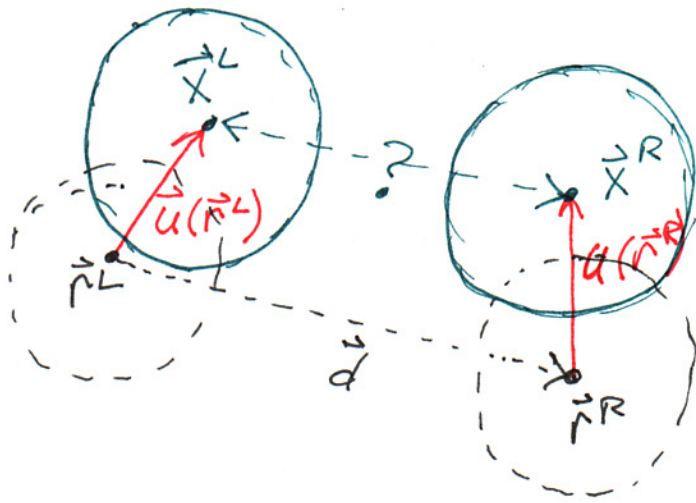
$$\text{Total Force} = 0 \quad \checkmark$$

$$\begin{aligned} \text{Total Torque} \\ = (\sigma_{yx} - \sigma_{xy}) \cdot L = 0 \end{aligned}$$

$$\Rightarrow \sigma_{xy} = \sigma_{yx}$$

The Strain Tensor

[Standard derivation in texts: this fancy one from Landau & Lifshitz ~~Elasticity~~ Theory of Elasticity p.1]



Two atoms, undistorted positions \vec{r}^L, \vec{r}^R , distorted positions \vec{x}^L, \vec{x}^R . Displacement field

$$\vec{u}(\vec{r}) = \vec{x}(\vec{r}) - \vec{r}$$

Original distance

$$d^z = (\vec{r}^R - \vec{r}^L)^2 \quad \text{Final distance}$$

Strain tensor tells how far the original distance vector \vec{d} will stretch.

$$\begin{aligned} (\vec{x}^R - \vec{x}^L)^2 &= [\vec{r}^R + \vec{u}(\vec{r}^R) - \vec{r}^L - \vec{u}(\vec{r}^L)]^2 \\ &= [d_i + u_i(\vec{r}^L + \vec{d}) - \bar{u}_i(\vec{r}^L)]^2 \quad \leftarrow \text{Summation convention } a_i^2 = a_i \cdot a_i = \sum_{i=1}^3 a_i^2 \\ &\approx [d_i + d_j \partial_j u_i]^2 \quad \leftarrow \vec{u}(\vec{x} + \vec{h}) - \vec{u}(\vec{x}) \approx (\vec{h} \cdot \nabla) \vec{u} \end{aligned}$$

Old distance

$$= d_i^2 + 2 d_i d_j \partial_j u_i + d_j \partial_j u_i d_k \partial_k u_i \quad \leftarrow \text{true for vectors } a \text{ \& } \text{well as functions}$$

$$\begin{aligned} d_i d_j \partial_j u_i &= d_i d_j \partial_i u_j \\ &= d_i d_j \left(\frac{\partial_j u_i + \partial_i u_j}{2} \right) \end{aligned}$$

Interchange dummy indices \underline{i} and \underline{k}

$$\begin{aligned}
 & (\text{New distance})^2 - (\text{Old distance})^2 \\
 &= 2 \, d_i d_j \left[\frac{\partial_i u_j + \partial_j u_i}{2} + \frac{\partial_i u_k \partial_j u_k}{2} \right] \\
 &= 2 \, d_i d_j \, \epsilon_{ij}
 \end{aligned}$$

Strain tensor

$$\epsilon_{ij} = \frac{1}{2} \left[\underbrace{\partial_i u_j + \partial_j u_i}_{\text{Symmetrized Gradient}} + \underbrace{\partial_i u_k \partial_j u_k}_{\text{Geometric Nonlinearity}} \right]$$

The geometric nonlinearity is usually ignored. Strains for most materials (not rubber) above 1% lead to permanent bending (plastic flow).

$$\epsilon_{ij} \sim 0.01 \Rightarrow \partial_i u_k \partial_j u_k \sim 10^{-4}, \text{ ignore.}$$

Except: Rubber

Large rotations (wires, thin sheets, grain boundaries)

The Tensor of Elasticity

Hooke's Law $F = KX$, $F/A = Y \frac{\Delta L}{L}$, ...

Force \leftrightarrow Stress tensor σ_{ij}

Stretch \leftrightarrow Strain tensor ϵ_{kl}

$$\sigma_{ij} = C_{ijkl} \epsilon_{kl}$$

Elastic Tensor

- How many elastic constants for an anisotropic material?

$i, j, k, l \in 1, 2, 3 \Rightarrow 81$ at most.

$\epsilon_{kl} = \epsilon_{lk} \rightarrow (C_{ijkl} + C_{ijlk})$ value irrelevant for σ

\rightarrow set $C_{ijkl} = C_{ijlk}$

$C_{ijkl} = C_{jike}$

Six choices of $\{ij\}$, six choices of $\{kl\}$

$\rightarrow 36$ at most.

Use formula for the energy, ~~Problem 7.6.2 &~~

Problem 7.4.1 \Rightarrow 21 possible independent elastic constants.