

Miller *et al.* monolithically fabricate a series of 2D electron device microstructures that enable them to determine the effects of constriction of the 5/2 FQH fluid. They find that constriction by a 0.5- μm -wide quantum point contact does indeed destroy the 5/2 FQH state. But, they also find that confinement by slightly larger constrictions, of just 0.8 and 1.2 μm wide, doesn't. This more encouraging result has two important implications for our ability to investigate non-abelian statistics experimentally.

First, the survival of the fragile 5/2 state within a single constriction could enable the fractional charge of the quasiparticles that populate it to be determined, and compared to the value of $e/4$ predicted for the Moore–Read state. In previous studies, the charge of the quasiparticles of the abelian 1/3 FQH state has been measured either by monitoring the one-by-one addition of quasiparticles into an anti-dot placed within a constriction⁵, or by analysing the power spectrum of the shot noise that arises from tunnelling of quasiparticles through the constriction^{6,7}. In principle, the same approaches should be applicable to the authors' 5/2 FQH system. Yet, the very

non-abelian nature of the quasiparticles of this system might yet throw up some surprises that would complicate the answer to the question of whether the Moore–Read formalism paints an accurate picture of the 5/2 FQH state.

But perhaps a more ambitious, and certainly more exciting, possibility offered by the present work is to build complex device structures with multiple constrictions. This would enable the construction of a quasiparticle Aharonov–Bohm interferometer (see Fig. 1), similar to that used to probe the abelian anyon statistics of the 1/3 FQH fluid⁸. Theoretical analysis^{9,10} of such an interferometer operating in a non-abelian 5/2 FQH regime predicts the emergence of unusual, history-dependent interference phenomena. One of the key challenges to building such a device will be to engineer constrictions that are narrow enough to transmit measurable quasiparticle tunnelling currents while maintaining a sufficiently small electron depletion to ensure that the 5/2 FQH fluid remains dominant. For the strong 1/3 FQH state it was adequate to limit the constriction electron density depletion to

7%, relative to the centre of island within the Aharonov–Bohm ring, to maintain the 1/3 fluid throughout the island⁸. But, in the case of the 5/2 state, because of the proximity and strength of competing nearby correlated electron states, it is likely that a much weaker depletion of around 1% would be needed to avoid the complications that a contribution by such states would cause. In the widest constrictions demonstrated by Miller *et al.*¹ the depletion is estimated to be around 16%, far short of this target. But if future work can bring this down to the level needed to study a pure non-abelian FQH fluid, the rewards in terms of new physics is likely to be great indeed.

References

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STATISTICAL MECHANICS

Crackling crossover

Magnetic domains in a thin film grow in a jerky manner as avalanches of spins flip their directions. At low temperatures, the measured distribution of avalanche sizes agrees with one theory; at high temperatures, with another.

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We are rapidly learning how to understand crackling noise. Many things crackle: paper when it is crumpled, faults when tectonic plates rub together and magnets as they change magnetization — the subject of the article by Ryu, Akinaga and Shin on page 547 of this issue¹. The response to a smooth, slow external force is a series of abrupt avalanches with a broad range of sizes. Thus earthquakes are the crackling response of the Earth as the continents drift.

Why do things crackle? The first hint comes from the simple power-law distribution of avalanche sizes. Most

earthquakes are small; there are many more magnitude-six earthquakes than magnitude seven, and very few of magnitude eight. Indeed, the probability $P(s)$ of an earthquake (or magnetic spin avalanche) of size s falls off as a power law, $s^{-\tau}$. The exponent τ is universal — meaning that it will be shared among a large family of materials and systems. It will typically differ between systems that are fundamentally different; for example, crumpled paper and earthquake faults respond in basically different ways to their external forces. The exponent will also depend on the dimensionality of the system; for example, thin two-dimensional films will be different from bulk three-dimensional magnets. But τ usually will be independent of the microscopic details of the materials (which enables theoretical models to describe real experiments accurately). Sometimes τ will be shared between

strikingly different systems — for example, magnets and fluids invading porous rock — that are in the same universality class.

Power laws provide a hint to understanding crackling noise because they suggest the existence of an emergent symmetry of the system: scale invariance. Scale invariance means that the system looks (statistically) the same when put under a magnifying glass. Figure 1 shows some of the avalanches in a simple model of a magnet. The wide variety of avalanche sizes and their characteristic fractal shapes are typical of scale-invariant systems. When magnified the small avalanches become medium sized, the medium sized become large, and one of the large ones might become the background 'infinite' avalanche — riddled with holes of all sizes formed by interior sub-avalanches. The domain walls observed by Ryu *et al.*¹ have this fractal scale-invariant symmetry too.

